

# Thermal Stresses of a Thick Circular Plate: Steady State Problem

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**ABSTRACT-** In this paper, an attempt has been made to study thermoelastic response of a thick circular plate occupying the space  $D: 0 \leq r \leq a, -h \leq z \leq h$ , with radiation type boundary conditions. We apply integral transform technique to find the thermoelastic solution.

**Keywords:** Thermo elastic Response, thick circular plate, Thermal Stresses.

## I. INTRODUCTION

Khobragade et al. [3 - 12] have derived temperature distribution, displacement function, thermal stresses and thermal deflection of a thick and thin circular plate. Further Khobragade et al. [13] have established displacement function, temperature distribution and stresses and deflection of a triangular plate.

This paper is concerned with transient thermoelastic problem of a thick circular plate occupying the space  $D: 0 \leq r \leq a, -h \leq z \leq h$ , due to heat generation with radiation type boundary conditions.

## II. STATEMENT OF THE PROBLEM

Consider thick circular plate of thickness  $2h$  occupying the space  $D: 0 \leq r \leq a, -h \leq z \leq h$ , the material is homogenous and isotropic. The differential equation governing the displacement potential function  $\phi(r, z, t)$  as Nowacki [2] is

$$\frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{\partial^2 \phi}{\partial z^2} = \left( \frac{1+\nu}{1-\nu} \right) \alpha_t T \quad (1)$$

Where  $\nu$  and  $\alpha_t$  are Poisson's ratio and linear coefficient of thermal expansion of the material of the plate and  $T$  is the temperature of the plate satisfying the differential equation as Noda [3] is

$$\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial z^2} + \frac{g(r, z)}{k} = 0 \quad (2)$$

The boundary conditions are

$$M_r(T, 0, 1, a) = g(z), \quad -h \leq z \leq h, \quad (3)$$

$$\left. \begin{aligned} M_z(T, 1, k_1, h) &= f_1(r) \\ M_z(T, 1, k_2, -h) &= f_2(r) \end{aligned} \right\}, \quad 0 \leq r \leq a, \quad (4)$$

Where  $k$  is thermal diffusivity of material of the plate. The displacement functions in the cylindrical coordinate system are represented by as Khobragade [4] are

$$u_r = \frac{\partial \phi}{\partial r} - \frac{\partial^2 L}{\partial r \partial z} \quad (5)$$

$$u_z = \frac{\partial \phi}{\partial z} + 2(1-\nu) \nabla^2 L - \frac{\partial^2 L}{\partial z^2} \quad (6)$$

Where  $L$  is the Love's function [14] and must satisfy

$$\nabla^2 \nabla^2 L = 0 \quad (7)$$

Where

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\partial^2}{\partial z^2}$$

The component of stresses are represented by the thermoelastic displacement potential  $\phi$  and Love's function  $L$  as Noda [3] are

$$\sigma_{rr} = 2G \left\{ \left( \frac{\partial^2 \phi}{\partial r^2} - \nabla^2 \phi \right) + \frac{\partial}{\partial z} \left( \nu \nabla^2 L - \frac{\partial^2 L}{\partial r^2} \right) \right\} \quad (8)$$

$$\sigma_{\theta\theta} = 2G \left\{ \left( \frac{1}{r} \frac{\partial \phi}{\partial r} - \nabla^2 \phi \right) + \frac{\partial}{\partial z} \left( \nu \nabla^2 L - \frac{1}{r} \frac{\partial^2 L}{\partial r^2} \right) \right\} \quad (9)$$

$$\sigma_{zz} = 2G \left\{ \left( \frac{\partial^2 \phi}{\partial z^2} - \nabla^2 \phi \right) + \frac{\partial}{\partial z} \left\{ \left( (2-\nu) \nabla^2 L - \frac{\partial^2 L}{\partial z^2} \right) \right\} \right\} \quad (10)$$

$$\sigma_{rz} = 2G \left\{ \frac{\partial^2 \phi}{\partial r \partial z} + \frac{\partial}{\partial r} \left\{ \left( (1-\nu) \nabla^2 L - \frac{\partial^2 L}{\partial z^2} \right) \right\} \right\} \quad (11)$$

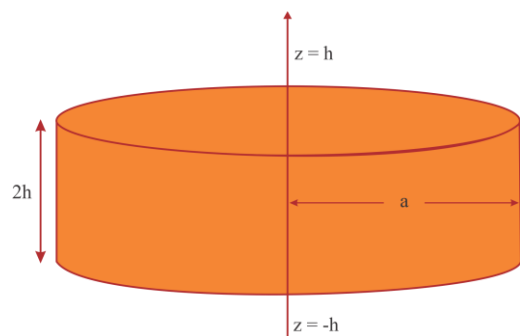


Fig. 1: Shows the geometry of the problem

For traction free surface stress function  $\sigma_z = \sigma_{r\theta} = 0$  at  $z = \pm h$  for thick plate

Equations (1) to (11) constitute the mathematical formulation of the problem under consideration.

## III. SOLUTION OF THE PROBLEM

Applying Marchi-Fasulo transform to the equation (2), we get

$$\frac{d^2\bar{T}}{dr^2} + \frac{1}{r} \frac{d\bar{T}}{dr} - \lambda_n^2 \bar{T} = \psi \quad (12)$$

$$\text{where } \psi = \frac{P_n(-h)}{k_2} f_2(r) - \frac{P_n(h)}{k_1} f_1(r) - \frac{\bar{g}}{k}$$

Equation (12) is a Bessel's equation whose solution is given by

$$\bar{T} = AI_0(\lambda_n r) + BK_0(\lambda_n r) + P.I. \quad (13)$$

Now, as  $r \rightarrow 0$ ,  $K_0 \rightarrow \infty$ , but  $\bar{T}$  is finite

$$\therefore B = 0.$$

Using (3), we get

$$A = \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)}$$

Substituting these values in equation (13) we get

$$\therefore \bar{T} = \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \quad (14)$$

Applying inverse Marchi-Fasulo transform we obtain

$$T = \sum_{n=1}^{\infty} \frac{P_n(z)}{\lambda_n} \left[ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \right] \quad (15)$$

Where,

$$P_n(z) = Q_n \cos(a_n z) - W_n \sin(a_n z),$$

$$Q_n = a_n(\alpha_1 + \alpha_2) \cos(a_n h) + (\beta_1 - \beta_2) \sin(a_n h)$$

$$W_n = (\beta_1 - \beta_2) \cos(a_n h) + a_n(\alpha_1 - \alpha_2) \sin(a_n h)$$

Equation (15) is the desired solution of the given problem.

Let us assume Love's function  $L$ , which satisfy condition (10) as

$$L = \sum_{n=1}^{\infty} \frac{P_n(z)}{\lambda_n} \left[ \frac{(\bar{f} - P.I. \Big|_{r=\xi})}{I_0(\lambda_n \xi)} I_0(\lambda_n r) \right] \quad (16)$$

Using (1) and (16), we get displacement potential  $\phi$  as

$$\therefore \phi = \frac{r^2 \alpha_t (1+\nu)}{4(1-\nu)} \sum \frac{P_n(z)}{\lambda_n} \times \left[ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \right] \quad (17)$$

#### IV. DETERMINATION OF DISPLACEMENT FUNCTION

Substituting equations (16) and (17) in equation (6), we get

$$u_r = \frac{\alpha_t}{4} \left( \frac{1+\nu}{1-\nu} \right) \sum_{n=1}^{\infty} \left\{ \frac{P_n(z)}{\lambda_n} \times \frac{(\bar{f} - P.I. \Big|_{r=\xi})}{I_0(\lambda_n \xi)} \right. \\ \left. \left( r^2 \lambda_n I_0'(\lambda_n r) + 2r I_0(\lambda_n r) \right) + r^2 \frac{\partial}{\partial r} (P.I.) + 2r (P.I.) \right\} \\ - \sum_{n=1}^{\infty} P_n'(z) \times \frac{(\bar{f} - P.I. \Big|_{r=\xi})}{I_0(\lambda_n \xi)} I_0'(\lambda_n r) \quad (18)$$

$$u_z = \frac{r^2 \alpha_t}{4} \left( \frac{1+\nu}{1-\nu} \right) \sum_{n=1}^{\infty} \frac{P_n'(z)}{\lambda_n} \left[ \frac{(\bar{f} - P.I. \Big|_{r=\xi})}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \right] \\ + 2(1-\nu) \sum_{n=1}^{\infty} \frac{\bar{f} - P.I. \Big|_{r=\xi}}{\lambda_n I_0(\lambda_n \xi)} \\ \left[ P_n(z) \lambda_n^2 I_0''(\lambda_n r) + \frac{P_n(z)}{r} \lambda_n I_0'(\lambda_n r) + P_n''(z) I_0(\lambda_n r) \right] \\ - \sum_{n=1}^{\infty} \frac{P_n'(z)}{\lambda_n} \times \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \times I_0(\lambda_n r) \quad (19)$$

Substituting equations (16) and (17) in equations (9) to (12), we obtain

$$\sigma_{rr} = \frac{-G\alpha_t}{2} \left( \frac{1+\nu}{1-\nu} \right) \\ \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \left\{ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \left[ P_n(z) (r \lambda_n I_0'(\lambda_n r) + 2I_0(\lambda_n r)) + r P_n''(z) r^2 (P.I.) \right] \right. \\ \left. + 2G \sum_{n=1}^{\infty} \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \right. \\ \left. \left\{ P_n'(z) [(v-1) \lambda_n^2 I_0''(\lambda_n r) + \frac{\lambda_n I_0'(\lambda_n r)}{r} + P_n'''(z) I_0(\lambda_n r)] \right\} \right. \quad (20)$$

$$\sigma_{\theta\theta} = \frac{-G\alpha_t}{2} \left( \frac{1+\nu}{1-\nu} \right) \sum_{n=1}^{\infty} \frac{1}{\lambda_n} \\ \left\{ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \left[ P_n(z) (r \lambda_n^2 I_0''(\lambda_n r) + 4r \lambda_n I_0'(\lambda_n r)) + P_n''(z) r^2 I_0(\lambda_n r) \right] \right. \\ \left. + 2G \sum_{n=1}^{\infty} \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \right. \\ \left. + P_n(z) \left[ r^2 \frac{\partial^2}{\partial r^2} (P.I.) + 4r \frac{\partial}{\partial r} (P.I.) + (P.I.) \right] + P_n''(z) r^2 (P.I.) \right\}$$

$$\left\{ P_n'(z) \left[ \left( \nu - \frac{1}{r} \right) \lambda_n^2 I_0''(\lambda_n r) + \frac{\nu}{r} I_0(\lambda_n r) + P_n''(z) \nu I_0(\lambda_n r) \right] \right\} \quad (21)$$

$$\sigma_{zz} = \frac{-G\alpha_t}{2} \left( \frac{1+\nu}{1-\nu} \right) \sum_{n=1}^{\infty} \frac{P_n(z)}{\lambda_n} \left\{ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \left[ r^2 \lambda_n^2 I_0''(\lambda_n r) + 5r \lambda_n I_0'(\lambda_n r) + 4I_0(\lambda_n r) \right] \right. \\ \left. r^2 \frac{\partial^2}{\partial r^2} (P.I.) + 5r \frac{\partial}{\partial r} (P.I.) + (P.I.) + 4(P.I.) \right\} \\ + 2G \sum_{n=1}^{\infty} \frac{\bar{f} - P.I. \Big|_{r=\xi}}{\lambda_n I_0(\lambda_n \xi)} \\ \left\{ P_n'(z) \left[ (2-\nu) \lambda_n^2 I_0''(\lambda_n r) + \left( \frac{2-\nu}{r} \right) \lambda_n I_0'(\lambda_n r) \right] \right. \\ \left. + P_n'''(z) (1-\nu) I_0(\lambda_n r) \right\} \quad (22)$$

$$\sigma_{rz} = \frac{G\alpha_t}{2} \left( \frac{1+\nu}{1-\nu} \right) \sum_{n=1}^{\infty} \frac{P_n'(z)}{\lambda_n} \left\{ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \left[ r^2 \lambda_n I_0'(\lambda_n r) + 2r I_0(\lambda_n r) \right] + r^2 \frac{\partial}{\partial r} (P.I.) + 2r (P.I.) \right\} \\ + 2G \sum_{n=1}^{\infty} \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} \\ \left\{ P_n(z) \lambda_n \left[ (1-\nu) \lambda_n I_0'''(\lambda_n r) + \frac{1-\nu}{r} I_0''(\lambda_n r) - \left( \frac{1-\nu}{r^2} \right) I_0'(\lambda_n r) \right] \right. \\ \left. + P_n'''(z) \nu I_0'(\lambda_n r) \right\} \quad (23)$$

### V. SPECIAL CASE

$$\text{Set } F(r, z) = z^2(1-r^2) \quad (24)$$

Applying Marchi-Fasulo transform, we obtain

$$\bar{F}(r, n) = (1-r^2) \Phi_n \left[ \frac{2h^2 \sin(a_n h)}{a_n} + \frac{4h \cos(a_n h)}{a_n^2} - \frac{4 \sin(a_n h)}{a_n^3} \right] \quad (25)$$

where

$$\Phi_n = a_n(\alpha_1 + \alpha_2) \cos(a_n h) + (\beta_1 - \beta_2) \sin(a_n h).$$

Using equation (25) in equation (15), one obtains

$$T = \sum_{n=1}^{\infty} \frac{P_n(z)}{\lambda_n} \left[ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \right] \quad (26)$$

### VI. NUMERICAL RESULTS

Set  $a = 2, \xi = 1.5, k = 15.9 \times 10^6, t = 1$  second in equation (26), we get

$$T = \sum_{n=1}^{\infty} \frac{P_n(z)}{\lambda_n} \left[ \frac{\bar{f} - P.I. \Big|_{r=\xi}}{I_0(\lambda_n \xi)} I_0(\lambda_n r) + P.I. \right] \quad (27)$$

### VII. CONCLUSION

In this paper, the temperature distribution, displacement and thermal stresses of a thick circular plate are investigated with known boundary conditions. Finite integral transform technique have been used to obtain numerical results. The results are obtained in terms of Bessel's function in the form of infinite series.

Any particular cases of special interest can be assigned to the parameters and functions in expressions. The results that are obtained can be useful to the design of structure or machines in engineering applications.

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