

Heat Transfer and Thermal Stresses Of Semi Infinite Rectangular Beam

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Abstract- This paper is concerned with transient thermoelastic problem in which we need to determine the temperature distribution, displacement function and thermal stresses of a semi-infinite rectangular beam when the boundary conditions are known. Integral transform techniques are used to obtain the solution of the problem.

Key Words: Semi-infinite rectangular beam, heat transfer thermo elastic problem, Integral transform.

I. INTRODUCTION

Khobragade et al. [2-7, 9] have investigated temperature distribution, displacement function, and stresses of a thin rectangular plate and Khobragade et al. [8] have established displacement function , temperature distribution and stresses of a semi-infinite rectangular beam.

In this paper, an attempt has been made to determine the temperature distribution, displacement function and thermal stresses of a semi-infinite rectangular beam occupying the region $D : -a \leq x \leq a ; -b \leq y \leq b, 0 \leq z \leq \infty$ with known boundary conditions. Here Marchi-Fasulo transforms and Fourier cosine transform techniques have been used to find the solution of the problem.

II. STATEMENT OF THE PROBLEM

Consider a thin rectangular plate occupying the space $D: a \leq x \leq a ; -b \leq y \leq b, 0 \leq z \leq \infty$. The displacement components u_x and u_y u_z in the x and y and z directions respectively as Tanigawa et al. [1] are

$$u_x = \int_{-a}^a \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial y^2} + \frac{\partial^2 U}{\partial z^2} - \nu \frac{\partial^2 U}{\partial x^2} \right) + \lambda T \right] dx \quad (1)$$

$$u_y = \int_{-b}^b \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial z^2} + \frac{\partial^2 U}{\partial x^2} - \nu \frac{\partial^2 U}{\partial y^2} \right) + \lambda T \right] dy \quad (2)$$

$$u_z = \int_0^\infty \left[\frac{1}{E} \left(\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} - \nu \frac{\partial^2 U}{\partial z^2} \right) + \lambda T \right] dz \quad (3)$$

where E, ν , and λ are the young's modulus, Poisson's ratio and the linear coefficient of the thermal expansion of the material of the beam respectively and $U(x,y,z,t)$ is the Airy's stress functions which satisfy the differential equation as Tanigawa et al. [1] is

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)^2 U(x, y, z, t) = -\lambda E \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)$$

$$\times T(x, y, z, t) \quad (4)$$

where $T(x,y,z,t)$ denotes the temperature of a rectangular beam satisfy the following differential equation as Tanigawa et al. [1] is

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} + \frac{g(x, y, z, t)}{k} = \frac{1}{\alpha} \frac{\partial T}{\partial t} \quad (5)$$

where k is the thermal conductivity and α is the thermal diffusivity of the material,

subject to initial condition

$$T(x, y, z, 0) = 0 \quad (6)$$

The boundary conditions are

$$\left[T(x, y, z, t) + k_1 \frac{\partial T(x, y, z, t)}{\partial x} \right]_{x=a} = f_1(y, z, t) \quad (7)$$

$$\left[T(x, y, z, t) + k_2 \frac{\partial T(x, y, z, t)}{\partial x} \right]_{x=-a} = f_2(y, z, t) \quad (8)$$

$$\left[T(x, y, z, t) + k_3 \frac{\partial T(x, y, z, t)}{\partial y} \right]_{y=b} = f_3(x, z, t) \quad (9)$$

$$\left[T(x, y, z, t) + k_4 \frac{\partial T(x, y, z, t)}{\partial y} \right]_{y=-b} = f_4(x, z, t) \quad (10)$$

$$T(x, y, z, t) \Big|_{z=0} = 0 \quad (11)$$

$$T(x, y, z, t) \Big|_{z=\infty} = h(x, y, t) \quad (12)$$

The stress components in terms of $U(x, y, z, t)$ Tanigawa et al. [1] are given by

$$\sigma_{xx} = \left[\frac{\partial^2 U}{\partial y^2} + \frac{\partial^2 U}{\partial z^2} \right] \quad (13)$$

$$\sigma_{yy} = \left[\frac{\partial^2 U}{\partial z^2} + \frac{\partial^2 U}{\partial x^2} \right] \quad (14)$$

$$\sigma_{zz} = \left[\frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} \right] \quad (15)$$

The equations (1) to (16) constitute the mathematical formulation of the problem under consideration.

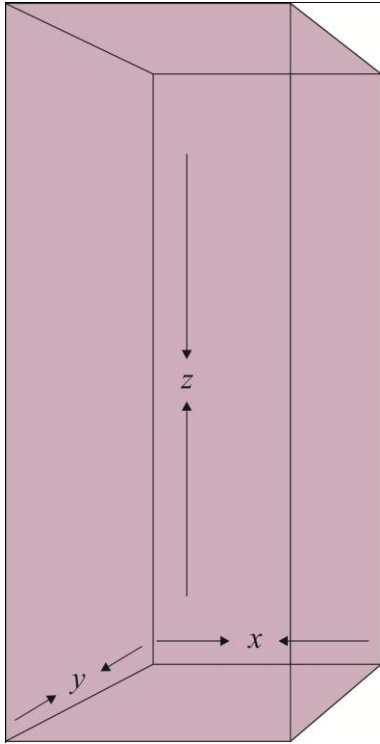


Fig 1: Geometry of the problem

III. SOLUTION OF THE PROBLEM

Applying finite Marchi-Fasulo transform over the variables x and y and finite Fourier sine transform over z, we get

$$\frac{d\bar{T}}{dt} + \alpha q^2 \bar{T} = \Psi \quad (16)$$

This is a linear equation whose solution is given by

$$\bar{T}(m, n, s, t) = \left(\int_0^t \Psi e^{-\alpha q^2(t-t')} dt' \right) \quad (17)$$

where, m, n, s are parameters of Marchi-Fasulo transform and sine transform respectively,

$$\Psi = \frac{P_m(a)}{k_1} \bar{f}_1^* - \frac{P_m(-a)}{k_2} \bar{f}_2^* + \frac{P_n(b)}{k_3} \bar{f}_3^* - \frac{P_n(-b)}{k_2} \bar{f}_4^* + \left[(-1)^{s+1} s \pi \bar{h}^* + \frac{g^*}{k} \right] \quad (18)$$

Now, applying inversion of Fourier sine transform and finite Marchi-Fasulo transform to the equation (18), one obtains the expression for temperature distribution as

$$T(x, y, z, t) = \left(\frac{2}{\pi} \right) \sum_{m,n=1}^{\infty} \left(\frac{P_m(x)P_n(y)}{\mu_m \zeta_n} \right) \times \int_0^{\infty} \left(\sin sz \int_0^t \Psi e^{-\alpha q^2(t-t')} dt' \right) ds \quad (19)$$

where,

$$q^2 = (\mu_m^2 + \zeta_n^2 + s^2 \pi^2)$$

Equation (19) is the required solution.

IV. AIRY'S STRESS FUNCTIONS

Substituting the value of temperature distribution T(x,y,z,t) from (20) in equation (4) one obtains

$$U(x, y, z, t) = \left(\frac{-2\lambda E}{\pi} \right) \sum_{m,n=1}^{\infty} \left(\frac{P_m(x)P_n(y)}{\mu_m \zeta_n} \right) \times \int_0^{\infty} B(t) \sin sz ds \quad (20)$$

Where

$$B(t) = \int_0^t \Psi e^{-\alpha q^2(t-t')} dt' \quad (21)$$

V. DISPLACEMENT COMPONENTS

Substituting the values of Airy's stress function from equation (22) in the equation (1) to (3), one obtains

$$u_x = \frac{-2\lambda}{\pi} \int_{-a}^a \int_{m,n=1}^{\infty} \left[\frac{1}{\mu_m \zeta_n} P_m(x)P_n''(y) - (s^2+1)P_m(x)P_n(y) - \nu P_m''(x)P_n(y) \sin(sz)B(t) \right] ds dx \quad (22)$$

$$u_y = \frac{-2\lambda}{\pi} \int_{-b}^b \int_{m,n=1}^{\infty} \left[\frac{1}{\mu_m \zeta_n} [-(s^2+1)P_m(x)P_n(y) + P_m^{(4,3pl)}(x)P_n(y) - \nu P_m(x)P_n''(y)] \sin(sz)B(t) \right] ds dy \quad (23)$$

$$u_z = \frac{-2\lambda}{\pi} \int_0^{\infty} \int_{m,n=1}^{\infty} \left[\frac{1}{\mu_m \zeta_n} P_m''(x)P_n(y) + P_m(x)P_n''(y) - \nu (s^2-1)P_l(x)P_m(y) \right] \sin(sz)B(t) ds dz \quad (24)$$

VI. DETERMINATION OF STRESS FUNCTION

Substituting the value of Airy's stress function U(x,y,z,t) from equation (22) in the equation (14) to (16) one obtain the stress functions as,

$$\sigma_{xx} = \frac{-2\lambda E}{\pi} \int_0^{\infty} \sum_{m,n=1}^{\infty} \frac{1}{\mu_m \zeta_n} \left[(P_m(x)P_n''(y) - P_m(x)P_n(y)s^2) \right] \times \sin(sz)B(t) ds \quad (26)$$

$$\sigma_{yy} = \frac{-2\lambda E}{\pi} \int_0^{\infty} \sum_{m,n=1}^{\infty} \frac{1}{\mu_m \zeta_n} \left[\frac{(P_m''(x)P_n(y)s^2 - P_m''(x)P_n''(y))}{\times \sin(sz)B(t)ds} \right] \quad (27)$$

$$\sigma_{zz} = \frac{-2\lambda E}{\pi} \int_0^{\infty} \sum_{m,n=1}^{\infty} \frac{1}{\mu_m \zeta_n} \left[\frac{(P_m''(x)P_n''(y) + P_m''(x)P_n''(y))}{\times \sin(sz)B(t)ds} \right] \quad (28)$$

VII. SPECIAL CASE

Set

$$h(x, y, t) = (x - a)^2(x + a)^2 \times (y - b)^2(y + b)^2(e^{-t})$$

$$\therefore \overset{=}{h}(n, m, s) = (z + e^{-z})(e^{y-t}) \times \left[\frac{a_n \cos^2(a_n a) - \cos(a_n a) \sin(a_n a)}{a_n^2} \right] \quad (29)$$

Substituting the above value in equation (19) one obtains

$$T(x, y, z, t) = \left(\frac{2}{\pi} \right) \sum_{m,n=1}^{\infty} \left(\frac{P_m(x)P_n(y)}{\mu_m \zeta_n} \right) \times \int_0^{\infty} \left(\sin sz \int_0^t \Psi e^{-\alpha q^2(t-t')} dt' \right) ds \quad (30)$$

VIII. NUMERICAL RESULTS

Set $a = 2, \alpha = 0.86, b = 3, t = 1$ sec in the equations (30) to obtain

$$T(x, y, z, t) = \left(\frac{2}{\pi} \right) \sum_{m,n=1}^{\infty} \left(\frac{P_m(x)P_n(y)}{\mu_m \zeta_n} \right) \times \left[\frac{a_n \cos^2(2a_n) - \cos(2a_n) \sin(2a_n)}{a_n^2} \right] \times \left[\frac{b_m \cos^2(3b_m) - \cos(3b_m) \sin(3b_m)}{b_m^2} \right] \times \int_0^{\infty} \left(\sin sz \int_0^1 \Psi e^{-t'} e^{-(0.86)q^2(1-t')} dt' \right) ds \quad (31)$$

IX. MATERIAL PROPERTIES

The numerical calculations has been carried out for an

Aluminum (pure) rectangular beam with the material properties as,

Density $\rho = 169$ lb/ft³

Specific heat = 0.208 Btu/lbOF

Thermal conductivity K = 117 Btu/(hr. ftOF)

Thermal diffusivity $\alpha = 3.33$ ft²/hr.

Poisson ratio $\nu = 0.35$

Coefficient of linear thermal expansion

$$\alpha_t = 12.84 \times 10^{-6} 1/F$$

Lame constant $\mu = 26.67$

Young's modulus of elasticity E = 70G Pa

X. DIMENSIONS

The constants associated with the numerical calculation are taken as

Length of rectangular beam x = 4ft

Breath of rectangular beam y = 3 ft

Height of rectangular beam z = 10³ft

XI. CONCLUSION

In this article, the temperature distribution, unknown temperature gradient, displacement function and thermal stresses of a semi-infinite rectangular beam have been obtained, when the boundary conditions are known with the aid of finite Marchi-Fasulo transform and semi-infinite Fourier cosine transform techniques. The results are obtain in the form of infinite series in terms of Bessel's function.

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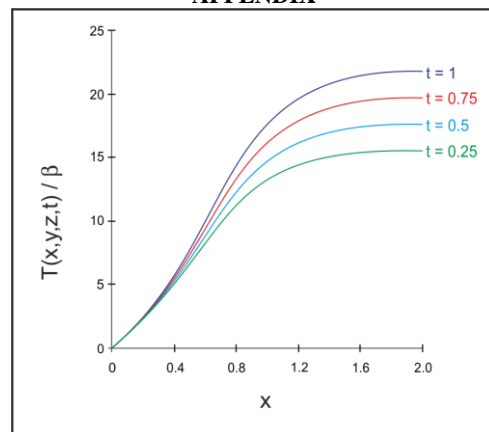


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APPENDIX



Graph 1: Graph of temperature distribution versus x